



# Habitable Zone Boundaries for Circumbinary Planets

Wolf Cukier<sup>1,2</sup>, Ravi kumar Kopparapu<sup>2,8</sup> , Stephen R. Kane<sup>3,8</sup> , William Welsh<sup>4</sup>, Eric Wolf<sup>5,8</sup>, Veselin Kostov<sup>2,6</sup>, and Jacob Haqq-Misra<sup>7,8</sup>

<sup>1</sup> Scarsdale High School, NY, USA

<sup>2</sup> NASA Goddard Space Flight Center, 8800 Greenbelt Road, Mail Stop 699.0, Building 34, Greenbelt, MD 20771, USA

<sup>3</sup> Department of Earth Sciences, University of California, Riverside, CA 92521, USA

<sup>4</sup> Department of Astronomy, San Diego State University, 5500 Campanile Drive, San Diego, CA 92182, USA

<sup>5</sup> Laboratory for Atmospheric and Space Physics, Department of Atmospheric and Oceanic Sciences, University of Colorado Boulder, Boulder, CO 80309, USA

<sup>6</sup> SETI Institute, 189 Bernardo Avenue, Suite 200, Mountain View, CA 94043, USA

<sup>7</sup> Blue Marble Space Institute of Science, 1001 4th Ave, Suite 3201, Seattle, WA 98154, USA

<sup>8</sup> NASA Astrobiology Institute's Virtual Planetary Laboratory, P.O. Box 351580, Seattle, WA 98195, USA

Received 2019 October 5; accepted 2019 October 23; published 2019 November 7

## Abstract

We use a one-dimensional (1D) cloud-free climate model to estimate habitable zone (HZ) boundaries for terrestrial planets of masses  $0.1 M_E$  and  $5 M_E$  around circumbinary stars of various spectral type combinations. Specifically, we consider binary systems with host spectral types F-F, F-G, F-K, F-M, G-G, G-K, G-M, K-K, K-M and M-M. Scaling the background  $N_2$  atmospheric pressure with the radius of the planet, we find that the inner edge of the HZ moves inwards toward the star for  $5 M_E$  compared to  $0.1 M_E$  planets for all spectral types. This is because the water-vapor column depth is smaller for larger planets and higher temperatures are needed before water vapor completely dominates the outgoing longwave radiation. The outer edge of the HZ changes little due to competing effects of the albedo and greenhouse effect. While these results are broadly consistent with the trend of single star HZ results for different mass planets, there are significant differences between single star and binary star systems for the inner edge of the HZ. Interesting combinations of stellar pairs from our 1D model results can be used to explore for in-depth climate studies with 3D climate models. We identify a common HZ stellar flux domain for all circumbinary spectral types.

*Key words:* astrobiology – binaries: close – planets and satellites: terrestrial planets

*Online material:* color figures

## 1. Introduction

Binary stars are ubiquitous in the galaxy, with nearly half of all Sun-like stars residing in binary (and higher multiple star) systems. Numerous studies in the past two decades have predicted that planets can form and sustain long-term stability around binary stellar systems (Pierens & Nelson 2007; Alexander 2012; Meschiari 2012; Paardekooper et al. 2012; Liu et al. 2013; Marzari et al. 2013; Mason et al. 2013, 2015; Georgakarakos & Egg 2015). The *Kepler* mission has detected several exoplanets in binary systems (Doyle et al. 2011; Orosz et al. 2012a, 2012b; Welsh et al. 2012, 2015; Kostov et al. 2013; Schwamb et al. 2013). Despite a strong observational bias against the discovery of such planets, at the time of writing there are six confirmed planets orbiting one member of a sub-20 au binary stellar system (i.e., circumprimary planets or S-type systems, Kley & Haghighipour 2014) and 12 confirmed planets orbiting within 3 au of both members of sub-astronomical unit binary star systems (circumbinary planets or P-type systems, e.g., Welsh et al. 2015; Kostov et al. 2016a). Based on known circumbinary systems, estimates suggest a 1%–10% occurrence

rate of Neptune- to Jupiter-sized planets (e.g., Armstrong et al. 2014; Welsh et al. 2015; Kostov et al. 2016b). Given the proximity to their host star, planets in binary systems experience the effects of two incident stellar fluxes. Almost half of known circumbinary planets reside in the habitable zone (HZ) (Doyle et al. 2011; Orosz et al. 2012a, 2012b; Welsh et al. 2015; Kostov et al. 2016b), as constrained by existing estimates of the HZ for binaries (Eggl et al. 2012, 2013; Haghighipour & Kaltenegger 2013; Kaltenegger & Haghighipour 2013; Kane & Hinkel 2013; Forgan 2016; Wang & Cuntz 2019).

Our goal in this study is to estimate the HZs around circumbinary terrestrial planets using one-dimensional (1D) climate models. We describe the methodology in Section 2, present the results of our analysis in Section 3, and provide a discussion of their implications in Section 4.

## 2. Model and Methods

We used the 1D radiative-convective, cloud-free climate model from the Kasting group, which has been updated recently (Kopparapu et al. 2013, 2014). Details of the model

are given in these papers and references there in. In order to simulate the stellar flux incident on a circumbinary planet, we follow the methodology of Kane & Hinkel (2013), where the combined stellar energy distribution (SED) from both the stars is used to estimate the equivalent effective temperature of a single energy source that would produce the same energy flux. We then ran the model 1000 times, each time at a different point in the orbit, where the stellar SED was combined. We performed these calculations for 10 cases: F-F, F-G, F-K, F-M, G-G, G-K, G-M, K-K, K-M and M-M. It is assumed that the planet is in a circular orbit around the stars, and the two stars are not orbiting each other. We used the BT-Settl grid of models<sup>9</sup> (Allard et al. 2003, 2007). Two end-member planetary masses are considered: 0.1 and 5  $M_E$ , to be consistent with previous studies (Kopparapu et al. 2014; Wang & Cuntz 2019). We scale the background  $N_2$  atmospheric pressure with the radius of the planet, which suggests that larger planets should have thicker atmospheres. The corresponding scaling is given in Kopparapu et al. (2014).

We followed the methodology from Kasting et al. (1993) and Kopparapu et al. (2013) to estimate the HZs. The inner edge of the HZ is calculated by increasing the surface temperature of a fully saturated Earth model from 220 K up to 2200 K. The effective solar flux  $S_{\text{eff}}$ , which is the value of solar constant required to maintain a given surface temperature, is calculated from the ratio between the net outgoing IR flux and the net incident solar flux, both evaluated at the top of the atmosphere. When  $S_{\text{eff}}$  asymptotes to a constant value, that is when the atmosphere is optically thick to the outgoing IR radiation, and the planet enters the runaway greenhouse regime. This is considered the inner edge of the HZ. The total flux incident at the top of the atmosphere is taken to be the present solar constant at Earth's orbit 1360 W m<sup>2</sup>. The outer edge of the HZ is calculated by fixing the surface temperature of an Earth-like planet with 1 bar  $N_2$  atmosphere, and the atmospheric  $CO_2$  partial pressure was varied from 1 to 35 bar. Due to competing effects of the outgoing IR and the incoming solar,  $S_{\text{eff}}$  experiences a minimum as a function of  $CO_2$  partial pressure. This minimum is where the “maximum” amount of warming can be achieved with  $CO_2$ , and this is the “maximum greenhouse” limit for the outer edge of the HZ.

### 3. Results

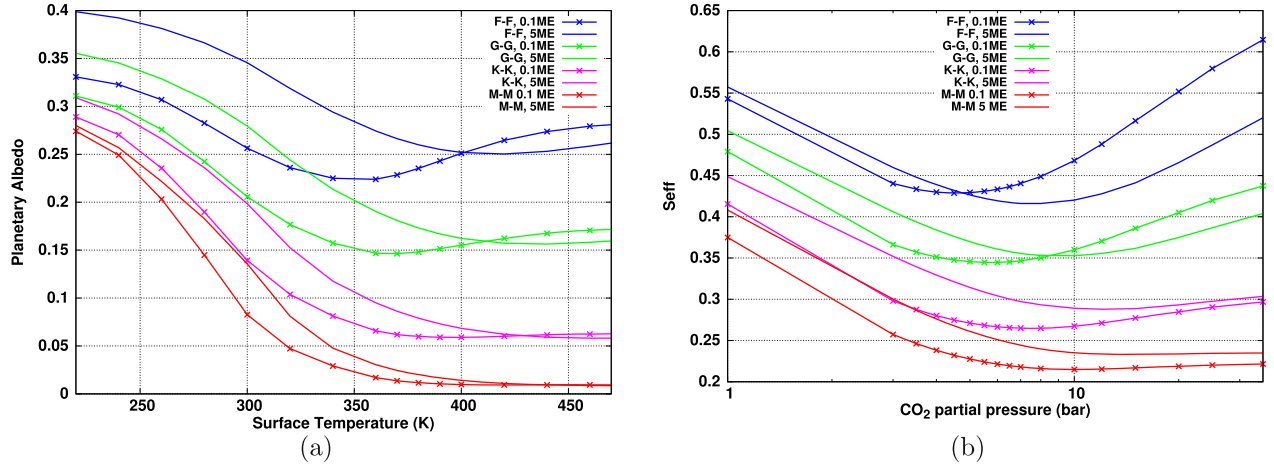
Figures 1(a) and (b) show the results for inner and outer HZ, respectively, of 0.1 and 5  $M_{\oplus}$  planets around binary stars of equal spectral types. Intermediate cases of mixed stellar spectral types (i.e., F-G, G-K, K-M etc.) fall within the regions of equal spectral types, and we do not show these results to maintain the clarity of our results shown in Figures 1(a) and (b).

The planetary albedo shown in Figure 1(a) is higher if the host binary is comprised of hotter stars, and lower if the binary has cooler stars. The reason is that the Rayleigh scattering cross section (inversely proportional to  $\lambda^4$ ) is on average higher for planets around an F star, as the star's Wien peak is bluer (shorter wavelength) compared to cooler stars. Furthermore,  $H_2O$  and  $CO_2$  have stronger absorption coefficients in the near-infrared (NIR) than in the visible, so the amount of starlight absorbed by the planet's atmosphere increases as the radiation is redder (as is the case for an M-M binary). Both effects are more pronounced when the atmosphere is dense and full of greenhouse gaseous absorbers. Hence, for a planet around M-M binary, the planetary albedo is significantly lower due to minimal Rayleigh scattering and high NIR absorption. Because we scaled the background  $N_2$  pressure with planetary mass, as was done in Kopparapu et al. (2014), there is a higher amount of non-condensable gas on 5  $M_E$  than on a 0.1  $M_E$ . This increases the Rayleigh scattering for a 5  $M_E$  (solid curves in Figure 1(a)) compared to a 0.1  $M_E$  (marker curve). This effect is more pronounced around host binary stars that have peak radiation more blueshifted (for ex: F-F), as can be seen in Figure 1(a). As the surface temperature increases, the planetary albedo decreases as a consequence of absorption of NIR solar radiation by  $H_2O$ . It then increases again and asymptotes at higher temperatures as Rayleigh scattering becomes dominant.

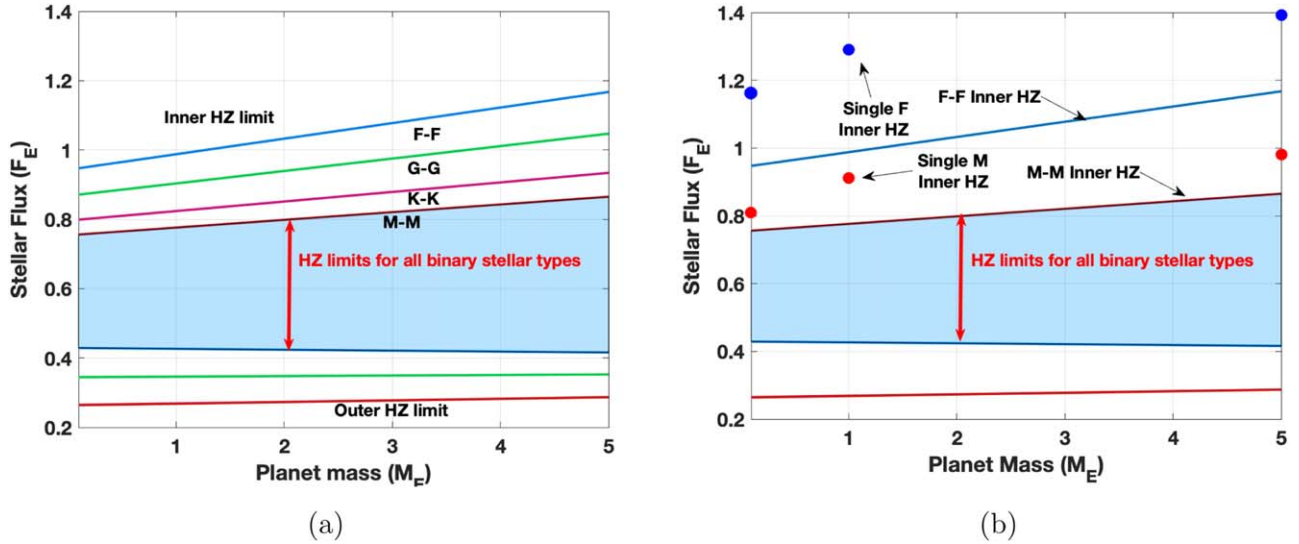
Figure 1(b) shows  $S_{\text{eff}}$  as function of partial pressure of  $CO_2$ , for our outer HZ calculations. The idea here is to estimate the *maximum* amount of  $CO_2$  needed to maintain a surface temperature of 273 K, and the corresponding energy balance needed to obtain this in the form of  $S_{\text{eff}}$ . As  $CO_2$  partial pressure increases,  $S_{\text{eff}}$  initially decreases because of the NIR absorption of  $CO_2$  greenhouse effect. However, beyond certain amount of  $CO_2$  in the atmosphere,  $CO_2$  ice clouds start forming and gradually increase the Rayleigh scattering (Kasting et al. 1993). This effect is not pronounced at first because the greenhouse effect of  $CO_2$  dominates at lower amounts of  $CO_2$ . Remember that  $S_{\text{eff}}$  is a ratio between the net outgoing IR flux and the net incident solar flux. As the amount of  $CO_2$  increases outgoing IR flux asymptotically approaches a constant value as the atmosphere becomes optically thick at all IR wavelengths. However, the net incident solar flux decreases monotonically with increases in  $CO_2$  partial pressure as a result of increased Rayleigh scattering. Hence,  $S_{\text{eff}}$  has a turnover, or a minimum at a corresponding partial pressure of  $CO_2$ . This is the *maximum* amount of  $CO_2$  that can provide a greenhouse warming, and hence the “maximum greenhouse limit.”

Comparing 0.1  $M_E$  and 5  $M_E$  planets for the outer HZ case (Figure 1(b)),  $S_{\text{eff}}$  is smaller for a lower mass planet, at low  $CO_2$  partial pressures. This is because the atmosphere of a 0.1  $M_E$  planet has a larger column depth, which increases the greenhouse effect and reduces the outgoing IR flux, decreasing  $S_{\text{eff}}$ . At the same time, this larger column depth also increases the planetary albedo at high  $CO_2$  partial pressures, *increasing*

<sup>9</sup> <http://perso.ens-lyon.fr/france.allard/>



**Figure 1.** (a): Top of the atmosphere planetary albedo as a function of surface temperature for a variety of spectral type configurations, and (b) effective stellar flux as a function of partial pressure of  $\text{CO}_2$ , for  $0.1 M_E$  (marker curve) and  $5 M_E$  (solid curve) planets around different host binary stars. (A color version of this figure is available in the online journal.)



**Figure 2.** HZ estimates from 0.1 to 5 Earth mass planets around all circumbinary stellar spectral types. The blue shaded region is the width of the HZ that is common across all stellar binary spectral types. The left panel (a) shows HZ limits for circumbinary star, while the right panel includes single star HZ stellar flux limits for comparison. The inner HZ limit for binary stars occurs at lower stellar fluxes (farther from the stars) because of the additional star contributing to the IR photons, and thus “raising” the near-IR flux incident on the planet, which increases the greenhouse warming. (A color version of this figure is available in the online journal.)

the  $S_{\text{eff}}$  for a  $0.1 M_E$  planet. A similar effect happens for the  $5 M_E$  mass planet, but here the atmosphere has a smaller column depth. Hence, it increases the outgoing IR flux (effectively “cooling” the planet) at low  $\text{CO}_2$  pressures, but at the same time Rayleigh scattering is weak at high  $\text{CO}_2$  pressures. The maximum greenhouse effect changes little because of these two competing effects, for a given host binary star system.

Figure 2 show the summary of the results from Figures 1(a) and (b), in terms of incident stellar flux on the vertical axis, and

planetary mass on the  $x$ -axis. Inner and outer HZs for F-F, G-G, K-K and M-M binary star types are shown in panels (a) and (b). The HZ limits for mixed combination of stellar spectral type binaries are a subset of these curves, and overlap within the ranges of the inner edge of the HZ for the F-F binary (top blue curve), and the outer edge of the M-M binary (bottom red curve). The blue shaded region represents the HZ for *all* different binary star and planetary mass scenarios. The inner edge of the HZ for a  $5 M_E$  occurs at relatively high stellar flux

compared to a  $0.1 M_E$  case, whereas the outer edge of the HZ changes very little. This result is consistent with the single star HZ results from Kopparapu et al. (2014), where they found that the HZ for a massive planet is larger compared to a smaller mass planet.

For comparison, we have also included HZ estimates for single stars (specifically, F and M spectral types) in Figure 2. The inner HZ limit in stellar flux for single stars is larger compared to binary stars because adding another star comparatively increases the number of photons available in the near-IR part of the combined SED. This enhances the greenhouse warming in comparison with the Rayleigh scattering, and one need not have to “push” the planet as close to the star (higher stellar flux) as for a single star to drive the planet into runaway greenhouse regime.

#### 4. Discussion

Our 1D model results assume a circular orbit for the planet around our binary star configurations. However, several exoplanets discovered around circumbinary stars have eccentric orbits. Previous studies have included the effect of eccentricity on Earth-like planets in the HZs of circumbinary planets (Eggl et al. 2012, 2013; Kane & Hinkel 2013; Mueller & Haghighipour 2014) and even eccentric host binary stars (Kley & Haghighipour 2014). These studies indicate that the orbit averaged flux incident on the planet varies significantly, depending upon the eccentricity. The time varying flux could induce changes in the climate of the planet, that could affect the habitability. While earlier studies have used the results of analytical or 1D climate model results to account for the impact on habitability, few 3D climate model results have been utilized to study the global dynamics of the planet (Popp & Eggl 2017). Future work using general circulation models (GCMs) is currently in progress from some of the co-authors of this study.

We have also compared our results with Eggl et al. (2013) and Kaltenegger & Haghighipour (2013). Comparing the data for a K-M stellar binary case (Haghighipour 2019, private communication), we find that Our data produced similar results to the Kaltenegger & Haghighipour (2013). However, Our data, at least for the inner edge of the HZ for a F-M binary, appears to differ by  $\sim 5\%$  to Eggl et al. (2013). This could be significant, depending upon how large is the planet. Kopparapu et al. (2014) found that the inner HZ for larger size planets move in by as much as 7%. The difference between Eggl et al. (2013) and our study likely to have arisen in the method of calculation that they propose (Eggl 2018).

There are substantial physical effects that are being ignored in this simple 1D model calculations (and any other result that is based on a non-higher dimensional model). For example, several 3D climate model results have shown that slow-synchronously rotating planets develop thick substellar clouds

due to weak Coriolis force. These substellar clouds increase the planetary albedo, potentially maintaining habitable conditions at higher stellar fluxes which otherwise would not be possible (Yang et al. 2013, 2014; Kopparapu et al. 2016; Way et al. 2016; Del Genio et al. 2018; Turbet et al. 2018). These general conclusions are more relevant at the inner edge of the HZ. At the outer edge, the models assume the regulation of the climate by the carbonate-silicate cycle. Recent calculations have suggested that planets in the outer regions of the HZ may be less likely to maintain stable, warm climates, but instead may oscillate between long, globally glaciated states and shorter periods of climatic warmth (Kadaya & Tajika 2014, 2015; Menou 2015; Haqq-Misra et al. 2016). Such conditions, similar to “Snowball Earth” episodes experienced on Earth, would be detrimental to the development of complex land life.  $\text{CO}_2$  sequestration beneath water ice, owing to the high density of  $\text{CO}_2$  ice ( $1.5 \text{ g cm}^{-3}$ ) compared to  $\text{H}_2\text{O}$  ice ( $1 \text{ g cm}^{-3}$ ), could potentially reduce or eliminate the deglaciation episodes (Turbet et al. 2017; Ramirez 2018). Limit cycles may also occur at higher stellar fluxes, near the inner edge of the HZ, at low outgassing rates (Paradise & Menou 2017). Both of these effects, which impact the inner and the outer HZs, need to be considered to improve upon the HZ estimates given in this work.

#### 5. Conclusions

We have estimated the HZs of 0.1 and 5 Earth mass planets around circumbinary stars of various stellar spectral types using a 1D radiative-convective climate model. We identify the width of the HZ that could be used for any spectral type of circumbinary star combination in circular orbit. We find that planetary albedo plays a major role in determining the inner edge of the HZ, while the competing effects of NIR absorption and the Rayleigh scattering of  $\text{CO}_2$  clouds at the outer edge make it less sensitive to the variations in these two parameters. However, a more detailed study using 3D climate model studies is needed, to properly consider the atmospheric circulation and the corresponding effect on the habitability of terrestrial planets around circumbinary stars. The value in this 1D study is identifying the interesting combinations of stellar pairs that would lead to interesting behavior in a GCM, as 1D studies are more suited for exploring parameter space. We thank the reviewer, Ramses Ramirez, for their constructive comments, which greatly improved the manuscript.

The authors gratefully acknowledge funding from NASA Habitable Worlds grant 80NSSC17K0741 for the support of this study. NASA affiliates acknowledge support from GSFC Sellers Exoplanet Environments Collaboration (SEEC), which is funded in part by the NASA Planetary Science Divisions Internal Scientist Funding Model.

## ORCID iDs

Ravi kumar Kopparapu  <https://orcid.org/0000-0002-5893-2471>

Stephen R. Kane  <https://orcid.org/0000-0002-7084-0529>

## References

- Alexander, R. 2012, *ApJL*, **757**, L29
- Allard, F., Allard, N. F., Homeier, D., et al. 2007, *A&A*, **474**, L21
- Allard, F., Guillot, T., Ludwig, H. G., et al. 2003, in Proc. IAU Symp. 211, Brown Dwarfs, 325, ed. E. Martín (San Francisco, CA: ASP)
- Armstrong, D. J., Osborn, H. P., & Brown, D. J. A. 2014, *MNRAS*, **444**, 2
- Del Genio, A. D., Way, M. J., Amundsen, D. S., et al. 2018, *AsBio*, **19**, 99
- Doyle, L. R., Carter, J. A., Fabrycky, D. C., et al. 2011, *Sci*, **333**, 6049
- Eggl, S. 2018, Handbook of Exoplanets, 61 (Berlin: Springer)
- Eggl, S., Pilat-Lohinger, E., Funk, B., et al. 2013, *MNRAS*, **428**, 3104
- Eggl, S., Pilat-Lohinger, E., Georgakarakos, N., et al. 2012, *ApJ*, **752**, 74
- Forgan, D. 2016, *MNRAS*, **463**, 2768
- Georgakarakos, N., & Egg, S. 2015, *ApJ*, **802**, 94
- Haghighipour, N., & Kaltenegger, L. 2013, *ApJ*, **777**, 166
- Haqq-Misra, J., Kopparapu, R. K., & Batalha, N. 2016, *ApJ*, **827**, 120
- Kadoya, S., & Tajika, E. 2014, *ApJ*, **790**, 107
- Kadoya, S., & Tajika, E. 2015, *ApJL*, **815**, L7
- Kaltenegger, L., & Haghighipour, N. 2013, *ApJ*, **777**, 2
- Kane, S. R., & Hinkel, N. 2013, *ApJ*, **762**, 7
- Kasting, J. F., Whitmire, D. P., & Reynolds, R. T. 1993, *Icar*, **101**, 108
- Kley, W., & Haghighipour, N. 2014, *A&A*, **564**, A72
- Kopparapu, R. K., Ramirez, R., Kasting, J. F., et al. 2013, *ApJ*, **765**, 131
- Kopparapu, R. K., Ramirez, R. M., SchottelKotte, J., et al. 2014, *ApJL*, **787**, L29
- Kopparapu, R. K., Wolf, E. T., Haqq-Misra, J., et al. 2016, *ApJ*, **819**, 1
- Kostov, V. B., McCullough, P. R., Hinse, T. C., et al. 2013, *ApJ*, **770**, 52
- Kostov, V. B., Orosz, J. A., Welsh, W. F., et al. 2016a, *ApJ*, **827**, 1
- Kostov, V. B., Moore, K., Tamayo, D., et al. 2016b, *ApJ*, **832**, 183
- Liu, H-G., Zhang, H., & Zhou, J-L. 2013, *ApJL*, **767**, L38
- Marzari, F., Thebault, P., Scholl, H., et al. 2013, *A&A*, **553**, A71
- Mason, P. A., Zuluaga, J. I., Clark, J. M., et al. 2013, *ApJL*, **774**, L26
- Mason, P. A., Zuluaga, J. I., Cuartas-Restrepo, P. A., et al. 2015, *IJAsB*, **14**, 3
- Menou, K. 2015, *E&PSL*, **429**, 20
- Meschiari, S. 2012, *ApJ*, **752**, 71
- Mueller, T. W. A., & Haghighipour, N. 2014, *ApJ*, **782**, 26
- Orosz, J. A., Welsh, W. F., Carter, J. A., et al. 2012a, *Sci*, **337**, 1511
- Orosz, J. A., Welsh, W. F., Carter, J. A., et al. 2012b, *ApJ*, **758**, 87
- Paardekooper, S-J., Leinhardt, Z. M., Thebault, P., et al. 2012, *ApJL*, **754**, L16
- Paradise, A., & Menou, K. 2017, *ApJ*, **848**, 33
- Pierens, A., & Nelson, R. P. 2007, *A&A*, **472**, 3
- Popp, M., & Eggl, S. 2017, *NatCo*, **8**, 14957
- Ramirez, R. 2018, *Geosc*, **8**, 280
- Schwamb, M. E., Orosz, J. A., & Carter, J. A. 2013, *ApJ*, **768**, 127
- Turbet, M., Bolmont, E., Leconte, J., et al. 2018, *A&A*, **612**, A86
- Turbet, M., Forget, F., Leconte, J., et al. 2017, *E&PSL*, **476**, 11
- Wang, Z., & Cuntz, M. 2019, *ApJ*, **873**, 113
- Way, M., Del Genio, A. D., Kiang, N. Y., et al. 2016, *GeoRL*, **43**, 8376
- Welsh, W. F., Orosz, J. A., & Carter, J. A. 2012, *Natur*, **7382**, 475
- Welsh, W. F., Orosz, J. A., & Short, D. R. 2015, *ApJ*, **809**, 26
- Yang, J., Cowan, N. B., & Abbot, D. S. 2013, *ApJL*, **771**, L45
- Yang, J., Gwenaël, B., Fabrycky, D., & Abbot, D. 2014, *ApJL*, **787**, L2